Structural Phase Transitions and Metallized Phenomena in Arsenic Telluride under High Pressure

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Table S1. Experimental lattice parameters and unit cell volumes of As_2Te_3 under high pressure

Phase	P (GPa)	a (Å)	b (Å)	c (Å)	ß (°)	$V(\mathring{\mathbf{A}}^3)$
α	0	14.3884(14)	4.0318(2)	9.9233(7)	94.96(1)	573.51(7)
α	0.58	14.2856(13)	4.0168(2)	9.8490(8)	95.11(1)	562.92(5)
α	1.10	14.1793(15)	4.0056(2)	9.7962(9)	95.22(1)	554.09(7)
α	2.10	14.0453(19)	3.9806(3)	9.6861(12)	95.47(1)	539.07(6)
α	3.09	13.9231(13)	3.9604(3)	9.6085(9)	95.58(1)	527.31(5)
α	4.15	13.8243(15)	3.9439(3)	9.5378(12)	95.50(1)	517.63(7)
α'	5.09	13.7653(13)	3.9297(3)	9.4835(12)	95.61(1)	510.28(5)
α'	6.16	13.6969(13)	3.9179(3)	9.441(1)	95.51(1)	504.30(5)
α'	7.16	13.6441(12)	3.9041(3)	9.3914(9)	95.48(1)	497.98(4)
α'	8.17	13.6069(13)	3.8934(3)	9.3250(8)	95.53(1)	491.72(5)
α'	9.18	13.5605(13)	3.8810(3)	9.2697(8)	95.51(1)	485.60(5)
α'	10.3	13.5167(11)	3.8713(4)	9.2214(8)	95.61(1)	480.23(5)
α'	11.3	13.4813(13)	3.8559(4)	9.1886(13)	95.52(1)	475.43(6)
α'	12.2	13.4495(17)	3.8459(4)	9.1515(14)	95.66(1)	471.06(9)
α''	13.2	13.455(2)	3.8401(7)	9.071(2)	95.50(1)	466.52(12)
α''	14.6	13.462(2)	3.8226(6)	9.0004(14)	95.34(1)	461.17(12)
α''	15.5	13.482(3)	3.8192(7)	8.9320(16)	95.21(2)	458.03(13)
α''	16.3	13.519(2)	3.8110(8)	8.8674(14)	95.02(2)	455.10(11)
α''	17.3	13.539(2)	3.8022(7)	8.7980(15)	94.95(2)	451.23(8)
$\alpha^{\prime\prime}$	18.4	13.574(4)	3.7918(10)	8.729(2)	94.69(3)	447.76(12)
γ	16.3	9.513(3)	6.9182(18)	9.678(4)	135.97(2)	442.81(18)
γ	17.3	9.482(3)	6.8865(14)	9.660(3)	135.94(2)	438.78(16)
γ	18.4	9.469(2)	6.8421(13)	9.652(3)	135.88(1)	435.45(15)
γ	19.4	9.445(2)	6.8255(11)	9.643(2)	135.91(1)	432.61(12)
γ	20.3	9.427(3)	6.7668(17)	9.624(3)	135.84(1)	427.74(14)
γ	21.5	9.408(4)	6.7282(15)	9.610(3)	135.91(2)	423.31(13)
γ	23.0	9.398(4)	6.6857(14)	9.599(3)	135.87(1)	420.05(12)
γ	24.6	9.359(3)	6.6388(13)	9.580(4)	135.80(1)	415.05(15)
γ	26.5	9.326(3)	6.5966(12)	9.561(2)	135.75(1)	410.49(6)
γ	29.0	9.298(2)	6.5291(10)	9.515(3)	135.66(1)	403.81(9)
γ	30.6	9.252(3)	6.4825(11)	9.506(2)	135.69(1)	398.25(4)
γ	33.4	9.225(2)	6.4127(10)	9.481(2)	135.68(1)	391.93(4)
γ	36.6	9.170(3)	6.3720(14)	9.4377(16)	135.68(1)	385.38(4)
γ	40.8	9.144(3)	6.2898(11)	9.4157(15)	135.72(1)	378.13(4)

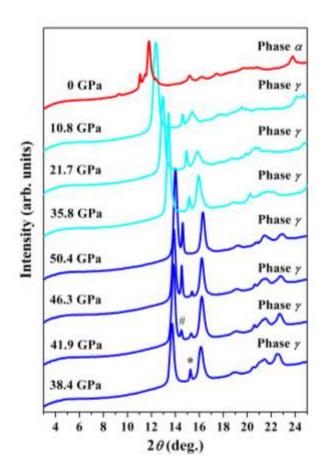


Figure S1. Selected angle-dispersive X-ray diffraction (AD-XRD) patterns of As₂Te₃ at room temperature in a loading process (blue lines) up to 50.4 GPa and a unloading process (cyan and red lines) down to 0 GPa ($\lambda = 0.6199 \text{ Å}$).

Notes:

In Figure S1, with increasing pressure, the asterisk (*) peak weakens and disappears at 50.4 GPa, while the intensity of number sign (#) peak increases upon compression. Due to the different width of these two peaks from those for phase γ , they do not belong to the high-pressure phases of As₂Te₃. At ambient pressure, there is not any impurity in the XRD pattern of the sample used in this work. In the unloading process, the number sign peak vanishes at 35.8 GPa and the asterisk peak reemerges and keeps down to at least 10.8 GPa. As shown in Figure S1, As₂Te₃ recovers to the original structure at pressure release, which shows that the phase transition could be reversible.

For the binary compounds composed of the V and VI group elements (As-Te, Bi-Te, Sb-Te, and Bi-Se materials), the chemical composition and crystal structures are very complex. S1-S4 According to Ref. S1a, it is easy for As₂Te₃ to decompose other compounds under high-pressure and high-temperature conditions. Therefore, part of As₂Te₃ may decompose into a new compound

indicated by the asterisk peak at high pressures and then this new material transforms into another phase (corresponding to the number sign peak) upon further compression. Because of lack of diffraction peaks and weak intensity of these new phases, it is difficult to obtain their chemical composition and crystal structures. Further high-pressure experiments would be performed at higher pressures to reveal the novel structural transitions in As_2Te_3 .

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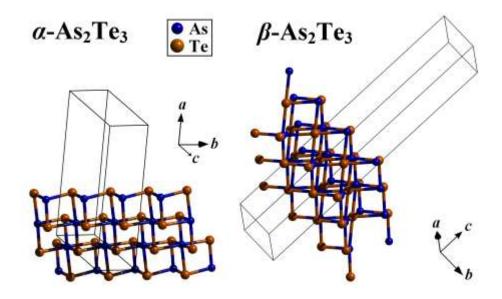


Figure S2. Schematic views of the $[As_4Te_6]_n$ ribbon structure in phase α and the Te(1)-As-Te(2)-As-Te(1) quintuple-layer structure in phase β of As_2Te_3 . The black bold lines denote the unit cells.

Notes:

In Figure S2, the $[As_4Te_6]_n$ ribbon in phase α and the Te(1)-As-Te(2)-As-Te(1) quintuple layer in phase β exhibit the 4 unit cells along the b-axis and the 3×3 unit cells parallel the ab-plane, respectively. The unit cell edges of these two phases are retained, in order to view distinctly the positions of ribbon or layer in the crystal lattice. Although the crystal structure of phase α of As_2Te_3 is different from that of phase β , the atomic arrangement in the $[As_4Te_6]_n$ ribbon in the former is similar to that in the quintuple layer in the latter. In the local ribbon or layer structures, As and Te atoms form two interpenetrating face-centered cubic (FCC) lattices, which is similar to that in the rock salt structure. Compared with the perfect rock salt structure, the distortions exist in the local structures of α - and β - As_2Te_3 . Although the $[As_4Te_6]_n$ ribbon in phase α is infinite along the b-axis, the atoms stack only one and two layers along the a- and c-axes, respectively. However, the Te(1)-As-Te(2)-As-Te(1) quintuple layer in phase β is an infinite two-dimensional structure parallel the ab-plane, in which the atoms stack five layers along the c-axis. Compared with the two-dimensional quintuple layer, the $[As_4Te_6]_n$ structure is only a one-dimensional ribbon composed of rock-salt-like As and Te atoms.